

Temperature dependence of structural and electronic properties of the spin-liquid candidate κ -(BEDT-TTF)₂Cu₂(CN)₃

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We investigate the effect that the temperature dependence of the crystal structure of a two dimensional organic charge-transfer salt has on the low-energy Hamiltonian representation of the electronic structure. For that, we determine the crystal structure of κ -(BEDT-TTF)₂Cu₂(CN)₃ for a series of temperatures between $T = 5$ K and 300 K by single crystal X-ray diffraction and analyze the evolution of the electronic structure with temperature by using density functional theory and tight binding methods. We find a considerable temperature dependence of the corresponding triangular lattice Hubbard Hamiltonian parameters. We conclude that even in the absence of change of symmetry, the temperature dependence of quantities like frustration and interaction strength can be significant and should be taken into account.

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I. INTRODUCTION

The two-dimensional organic charge-transfer salts based on bis(ethylenedithio)tetrathiafulvalene (BEDT-TTF or even shorter ET) molecules in a κ -type lattice arrangement have been intensively studied over the past thirty years due to their complex interplay between electron correlation and the effects of low dimensionality and spin frustration.¹ In particular, the discovery of spin-liquid behavior in κ -(BEDT-TTF)₂Cu₂(CN)₃² has fascinated experimentalists and theorists alike. Issues of current interest concern the nature of the low-temperature spin-liquid realized in this material^{3–5} and the various anomalies observed upon approaching the spin-liquid state from high temperatures. These anomalies include drastic changes in the ¹H-NMR relaxation rate around 200-150 K⁶, the thermopower at 150 K⁷, relaxor-type ferroelectricity around 60 K⁸ and a mysterious phase-transition anomaly at 6 K⁹. The latter feature, which manifests itself in anomalies in thermodynamic^{2,3} and transport⁴ quantities, is accompanied by pronounced lattice effects.⁹ Various scenarios have been suggested for the 6 K anomaly including a crossover from a thermally to a quantum disordered state³, an instability of the quantum spin-liquid^{3,10–16} or a distinct type of charge ordering.¹⁷ Theoretically, the spin-liquid properties have been investigated on the basis of the anisotropic triangular-lattice Hubbard Hamiltonian.^{12,18–21} The parameters t , t' and U of this Hamiltonian have been determined with semi-empirical⁷ as well as first principles methods^{22,23} based on the experimental structure at room temperature. Missing in such investigations is, however, the consideration of a possible temperature dependence of the model parameters.

In this work, we find by a combination of single crystal X-ray diffraction at various temperatures and density

functional theory calculations that even in the absence of structural phase transitions, the temperature dependence of the structural parameters is significant enough to influence the electronic behavior and the determination of the Hamiltonian model parameters in κ -(BEDT-TTF)₂Cu₂(CN)₃. We suggest that this has subtle effects on the degree of frustration and interaction strength.

In 1991, the crystal structure of κ -(BEDT-TTF)₂Cu₂(CN)₃ was first reported,²⁴ and the unit cell parameters confirmed,²⁵ on twinned crystals at room temperature. The crystal structure was redetermined in 1993, also on a twinned crystal at room temperature.²⁶ The crystal structure of the very similar κ' -(BEDT-TTF)₂Cu₂(CN)₃ structure, which is reported to be an ambient pressure superconductor, was described in 1992.^{27,28} The crystal structure of the κ phase was subsequently redetermined in 1997,²⁹ and the unit cell redetermined in 2001, both at room temperature.³⁰ Already in 1991, the unit cell parameters were reported at 300 and 30 K, with $a_{(30\text{ K})}/a_{(300\text{ K})} = 0.9964$, $b_{(30\text{ K})}/b_{(300\text{ K})} = 0.9932$, and $c_{(30\text{ K})}/c_{(300\text{ K})} = 0.9900$, and $\beta_{(30\text{ K})}/\beta_{(300\text{ K})} = 1.0146$.³¹ Our comparable contraction values between 300 and 20 K are $a_{(20\text{ K})}/a_{(300\text{ K})} = 0.9988$, $b_{(20\text{ K})}/b_{(300\text{ K})} = 0.9958$, and $c_{(20\text{ K})}/c_{(300\text{ K})} = 0.9905$, and $\beta_{(20\text{ K})}/\beta_{(300\text{ K})} = 1.0148$ and agree well with the previous results. In addition, we find evidence for an ordering of the ethylene groups in a staggered conformation between 200-150 K. Despite the current interest in this material as a spin-liquid candidate, no low-temperature structural determinations have yet been reported. Herein, we present a detailed characterization of the crystal structure as function of temperature determined on a single crystal with agreement factors better than any previously reported.

II. SINGLE CRYSTAL X-RAY DIFFRACTION

A black, plate-like crystal of κ -(BEDT-TTF) $_2$ Cu $_2$ (CN) $_3$ with dimension $0.4 \times 0.4 \times 0.4$ mm 3 was placed onto the tip of a glass fiber and mounted on a Bruker APEX II 3-circle diffractometer equipped with an APEX II detector. Temperature control in the 100-300 K region was provided by an Oxford Cryostream 700 Plus Cooler while below 100 K it was provided by a Cryocool-LHE cryogenic system (Cryo Industries of America). The sample temperature below 100 K was confirmed by installing a Cernox thermometer (Lakeshore) in the immediate vicinity of the crystal and stabilizing the temperature immediately prior to data collection. In order to reduce adverse thermal effects the thermometer wires (twisted manganin wires, 0.5 mm diameter, supplied by Lakeshore) were anchored on the surfaces of the cryostat exposed to 4 He stream. The data were collected using MoK $_{\alpha}$ radiation ($\lambda = 0.71073$ Å) with a detector distance of 50 mm. Unit cell parameters were determined upon cooling in 10 K increments between 100 and 290 K, with a frame exposure time of 10 seconds. Full data sets for structural analysis were collected at temperatures of 5, 20, 100, 150, 200, 250 and 300 K. The uncertainties in the temperature determination are typically ± 0.2 K down to 100 K, while for the 20 K and 5 K data points an error bar of ± 1 K has to be accepted. These data collections nominally covered over a hemisphere of reciprocal space by a combination of three sets of exposures. Data to a resolution of 0.68 Å were considered in the reduction. The raw intensity data were corrected for absorption (SADABS³²). The structure was solved and refined using SHELXTL.³³ A direct-method solution was calculated, which provided most of atomic positions from the E-map. Full-matrix least squares / difference Fourier cycles were performed, which located the remaining atoms. All non-hydrogen atoms were refined with anisotropic displacement parameters. The hydrogen atoms were placed in ideal positions and refined as riding atoms with relative isotropic displacement parameters.

As a representative example, in Fig. 1 we show the structure of κ -(BEDT-TTF) $_2$ Cu $_2$ (CN) $_3$ at a temperature of $T = 5$ K. In the bc plane, the BEDT-TTF molecules form their typical κ -type arrangement. One of the cyanide (CN $^-$) groups resides on an inversion center, thus requiring it to be disordered with a 50% carbon and 50% nitrogen distribution on these two atomic positions²⁴. Analysis of the low-temperature structural data (see Tab. I) indicates that overall, a , b and c -axes decrease with temperature while the β angle increases monotonically upon cooling. We find that the orthogonal projection of the a -axis, $a_{\perp} = a \sin \beta$, has the greatest relative contraction with temperature. One ethylene group is disordered at room temperature with a staggered conformation 77% of the time. As the temperature is lowered, the ethylene group remains partially disordered down to a temperature of 200 K and is fully ordered in a staggered

conformation at 150 K. The anion layer becomes slightly more buckled at low temperature: at room temperature, the Cu1 atom lies 0.050 Å out of the C11-C12-N12 plane, increasing to 0.072 Å at low temperature. For the purpose of performing density functional theory calculations, the original symmetry of $P2_1/c$ has been lowered to $P2_1$ by choosing one of two possible orientations of the CN $^-$ group in the inversion center.

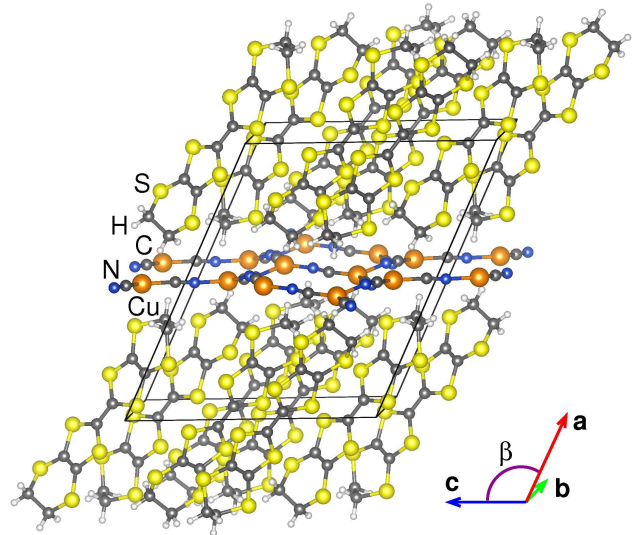


FIG. 1: Structure of κ -(BEDT-TTF) $_2$ Cu $_2$ (CN) $_3$ at $T = 5$ K. Note that the disordered CN $^-$ group in the inversion center is only shown in one conformation.

III. TEMPERATURE-DEPENDENT STRUCTURAL PARAMETERS

In Fig. 2, we show the evolution of the lattice parameters with temperature. Over the large investigated temperature range from $T = 300$ K down to $T = 5$ K, the volume is monotonously decreasing with temperature (see Fig. 2 (a)). The monoclinic angle β between \mathbf{a} and \mathbf{c} lattice vectors (see Fig. 2 (b)) first rapidly increases upon cooling down to a temperature of $T = 200$ K, below which it increases more gradually. In Fig. 2 (c), the lattice parameters are displayed as symbols, normalized by their values at $T = 20$ K (see Tab. I). We also include the relative b and c lattice constants along the two principal axes as obtained by thermal expansion measurements in Refs. 9 and 34. The out-of-plane expansivity data shown there were taken in a direction perpendicular to the bc plane. They are shown together with the corresponding quantity from the X-ray diffraction measurement, $a_{\perp} = a \sin \beta$.

	300 K	250 K	200 K	150 K	100 K	20 K	5 K
a (Å)	16.0919(3)	16.0848(3)	16.0781(3)	16.0703(3)	16.0746(6)	16.072(4)	16.062(3)
b (Å)	8.5722(2)	8.5749(1)	8.5737(1)	8.5664(2)	8.5593(3)	8.536(2)	8.544(2)
c (Å)	13.3889(2)	13.3373(2)	13.2964(2)	13.2698(3)	13.2678(5)	13.262(3)	13.271(2)
β (°)	113.406(1)	113.853(1)	114.273(1)	114.609(1)	14.852(1)	115.088(3)	115.093(2)
V (Å ³)	1694.93(6)	1682.43(4)	1670.86(4)	1660.72(6)	1656.51(1)	1647.8(6)	1649.3(5)
ρ (g/cm ⁻³)	1.909	1.923	1.937	1.949	1.954	1.964	1.962
μ (mm ⁻¹)	2.266	2.283	2.299	2.313	2.319	2.331	2.329
GoF	1.057	1.053	1.091	1.157	1.293	1.261	1.109
R	0.0311	0.0269	0.0238	0.0225	0.0220	0.0182	0.0204
R_w	0.0838	0.0735	0.0654	0.0627	0.0629	0.0503	0.0515
staggered (%)	77	86	93	100	100	100	100
$d_{\text{Cu-NNC}}$ (Å)	0.050	0.056	0.064	0.069	0.073	0.072	0.072
ET tilt angle ϑ	66.55	66.25	66.00	65.79	65.64	65.53	65.52
$d_{\text{intradimer}}$	3.558	3.538	3.518	3.500	3.488	3.470	3.473

TABLE I: Crystal data and structure refinement of κ -(BEDT-TTF)₂Cu₂(CN)₃. Formula = C₂₃H₁₆Cu₂N₃S₁₆; formula weight $M_W = 974.43$, monoclinic, wavelength $\lambda = 0.71073$ Å, effective number of electrons in the crystal unit cell contributing to $F(000) = 978$, space group $P2_1/c$, $Z = 2$. Residual factor for the reflections $R_1 = \sum ||F_o| - |F_c|| / \sum |F_o|$; weighted residual factors $wR_2 = [\sum w(F_o^2 - F_c^2)^2 / \sum w(F_o^2)^2]^{1/2}$; $[I > 2\sigma(I)]$; least-squares goodness-of-fit parameter GoF = $[\sum w(F_o^2 - F_c^2)^2 / (N_d - N_p)]^{1/2}$. Staggered (%) is the percentage of ethylene groups that are in the staggered conformation. $d_{\text{Cu-NNC}}$ (Å) is the distance that the Cu(I) ion is out of the plane defined by the coordinated N, N and C atoms. The ET tilt angle ϑ is measured against the bc plane. $d_{\text{intradimer}}$ is the orthogonal distance between BEDT-TTF dimers.

IV. TEMPERATURE-DEPENDENT ELECTRONIC STRUCTURE

In the following, we determine the electronic properties for the resolved crystal structures at different temperatures by employing the all-electron full-potential local orbitals (FPLO)³⁵ basis. We perform all calculations on a $(6 \times 6 \times 6)$ k mesh with a generalized gradient approximation functional.³⁶ In Fig. 3 we present the electronic band structures for the various crystal structures. In the calculation, we used the staggered (majority) conformation of the BEDT-TTF molecules at all temperatures. The changes as function of temperature for the two bands at the Fermi level, corresponding to the antibonding combinations of the BEDT-TTF highest occupied molecular orbital (HOMO) levels, are relatively small. On the other hand, the occupied bands down to -0.7 eV below the Fermi level show a significant dependence on temperature. These bands derive from the bonding combination of BEDT-TTF HOMO levels and from the $[\text{Cu}_2(\text{CN})_3]^-$ anion layer. Overall, these bands show a bandwidth that decreases with increasing temperature. This can be explained by the volume increase as function of temperature (see Fig. 2 (a)).

Further analysis of the electronic structure requires a reliable identification of the bands deriving from the BEDT-TTF molecules. For that purpose, band weights have been calculated for all structures and added up for all atoms corresponding to the BEDT-TTF cation layers and to the $[\text{Cu}_2(\text{CN})_3]^-$ anion layers, respectively. In Fig. 4, blue circles and orange triangles stand for a

predominance of BEDT-TTF weight and $[\text{Cu}_2(\text{CN})_3]^-$ weight, respectively. This identification allows us to fit the BEDT-TTF derived bands to a tight binding Hamiltonian

$$H_{\text{TB}} = \sum_{ij,\sigma} t_{i-j} (c_{i\sigma}^\dagger c_{j\sigma} + \text{H.c.}), \quad (1)$$

where $c_{i\sigma}^\dagger$ ($c_{i\sigma}$) create (annihilate) an electron with spin σ at site i ; the sites correspond to the positions of the BEDT-TTF molecules, shown as balls in Fig. 5. While a good overall fit of the band structure can be achieved by including six neighbor BEDT-TTF molecule distances up to $d = 9.4$ Å, a very good fit describing also the small dispersion along the $B - \Gamma$ direction, as shown by lines in Fig. 4, requires fourteen neighbor distances up to $d = 14.5$ Å.

We now proceed to analyze the temperature dependence of the tight binding parameters corresponding to the network of BEDT-TTF dimers that are highlighted in Fig. 5. These dimers form a triangular lattice, with the hopping parameter t' connecting dimers to chains along the b direction and the hopping parameter t forming the 2D connections in c direction. We are interested in estimating the parameters of the Hubbard Hamiltonian for the anisotropic triangular lattice

$$H = \sum_{\langle ij \rangle, \sigma} t (c_{i\sigma}^\dagger c_{j\sigma} + \text{H.c.}) + \sum_{[ij], \sigma} t' (c_{i\sigma}^\dagger c_{j\sigma} + \text{H.c.}) \\ + U \sum_i \left(n_{i\uparrow} - \frac{1}{2} \right) \left(n_{i\downarrow} - \frac{1}{2} \right) \quad (2)$$

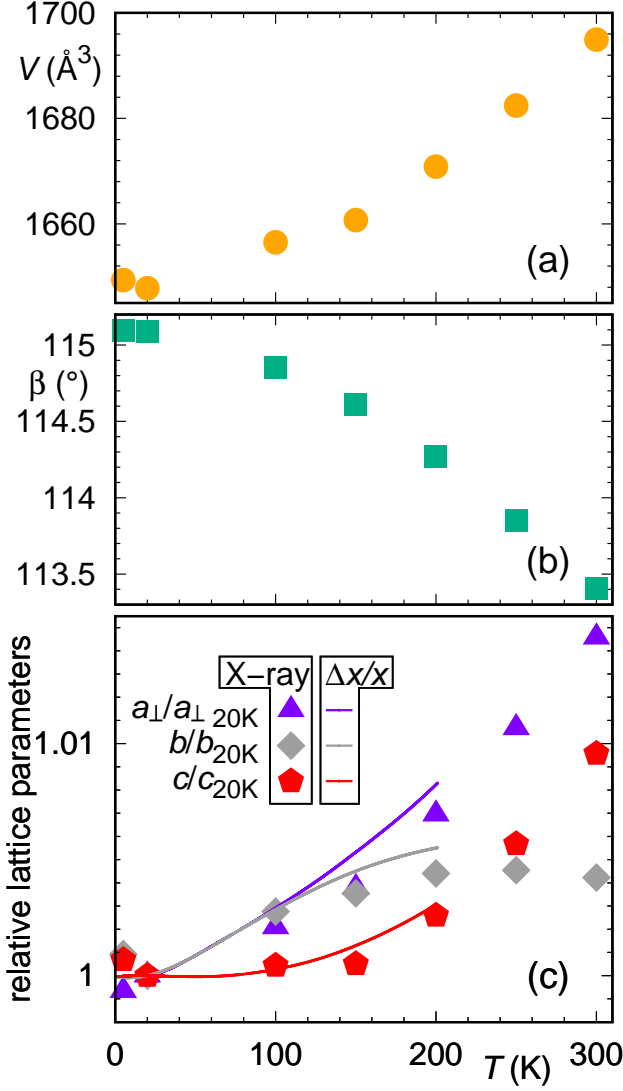


FIG. 2: Structural parameters of κ -(BEDT-TTF) $_2$ Cu $_2$ (CN) $_3$ between $T = 5$ K and $T = 300$ K. (a) and (b) show the volume and monoclinic angle, respectively. In (c) relative lattice parameters are given with the $T = 20$ K structure as reference. $a_{\perp} = a \sin \beta$. Symbols refer to the new data from X-ray diffraction (this work) while lines are measured thermal expansion data from Ref. 9. The experimental error bar is comparable to the size of the symbols.

where $\langle ij \rangle$ and $[ij]$ indicate summations over nearest and next-nearest neighbors, respectively. t and t' can be obtained from the molecular overlap integrals t_2 to t_4 by considering the geometrical formulas

$$t \approx \frac{t_2 + t_4}{2}, \quad t' \approx \frac{t_3}{2}. \quad (3)$$

For the definition of the overlap integrals $t_n = t_{i-j}$, see Ref. 22. Note that inclusion of longer-range hopping into these formulas (for example including t_5 into t') has no influence on the results reported in the following.

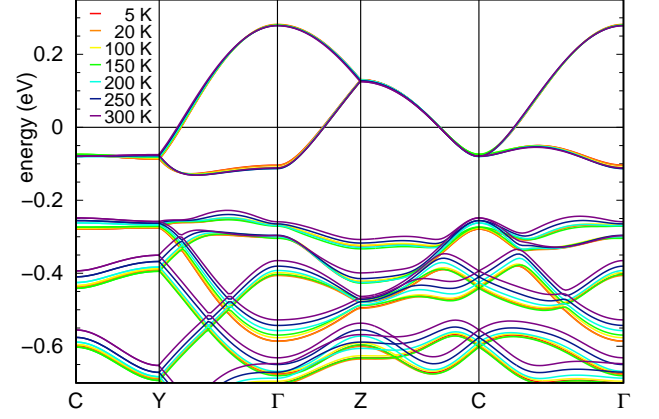


FIG. 3: Band structures of κ -(BEDT-TTF) $_2$ Cu $_2$ (CN) $_3$ from $T = 5$ K to $T = 300$ K.

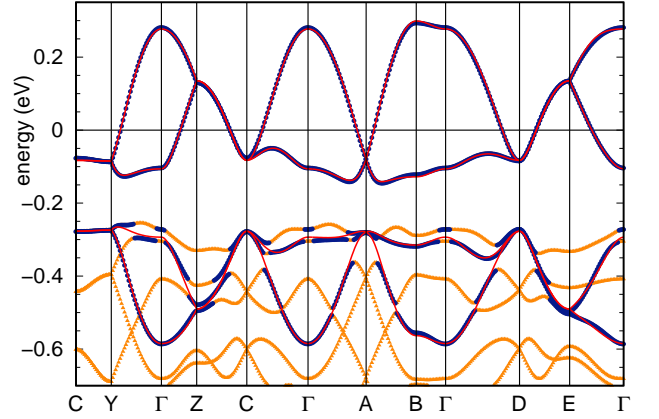


FIG. 4: Band structure of κ -(BEDT-TTF) $_2$ Cu $_2$ (CN) $_3$ at $T = 20$ K. Blue circles (orange triangles) indicate bands with a majority of BEDT-TTF and $[\text{Cu}_2(\text{CN})_3]^-$ character, respectively. The TB fit is shown with lines.

V. DISCUSSION

Fig. 6 summarizes our findings from the tight binding analysis. The nearest-neighbor hopping parameters t , forming a square lattice, increase upon cooling down to a temperature of $T = 200$ K, then decrease again (see Fig. 6 (a)). The frustrating hopping parameters t' show the opposite behavior as a function of temperature, decreasing upon cooling down to a temperature of $T = 150$ K, then increasing. Interestingly, these two effects enhance each other when we consider their ratio t'/t as shown in Fig. 6 (b)). t'/t which quantifies the degree of frustration in the system decreases from $t'/t = 0.82$ at $T = 300$ K to $t'/t = 0.80$ at $T = 150$ K, then increases to a maximum value of $t'/t = 0.86$ at $T = 5$ K. A rough estimate for the Coulomb interaction strength U can also be extracted from the dimer approximation $U \approx 2t_1$ where t_1 is the BEDT-TTF intradimer hopping integral (see Fig. 6 (c)). We find that the measure of the interaction

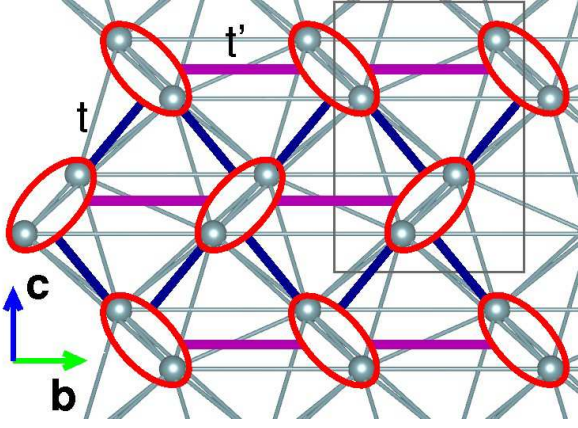


FIG. 5: Network formed by the BEDT-TTF molecules in the bc plane of κ -(BEDT-TTF) $_2$ Cu $_2$ (CN) $_3$. Each grey ball represents the center of gravity of a BEDT-TTF molecule, with grey bonds indicating the intermolecular distances that were taken into account for the tight binding fit. BEDT-TTF dimers are highlighted by ellipses, and the paths forming the triangular lattice paths are shown with bold lines. The unit cell is marked as a rectangle.

strength U/t estimated in this way monotonously falls by 7% as the temperature is increased from $T = 5$ K to $T = 300$ K.

In order to rationalize the observed temperature dependence of the Hamiltonian parameters in (2) we analyze the crystal structure in more detail. For that purpose, we determined the orientation of the BEDT-TTF molecules in space by finding the plane of the TTF part of the molecule and measuring its angle with respect to the bc plane, cf. inset of Fig. 7. This yields the inclination of the BEDT-TTF molecules against the anion layer shown as squares in Fig. 7 and the intradimer distance shown as circles (see also Table I). Both quantities show a nearly monotonous increase over the studied temperature range. The decreasing intradimer distance explains the increase of the intradimer hopping integral t_1 with decreasing temperature. Apparently, the nonmonotonous evolution of the overlap integrals t, t' , especially the distinct extrema in both quantities around 150 K to 200 K, has to be of different origin.

The overall trend can be understood by considering the temperature dependence of the lattice parameters; from Fig. 5 it is clear that changes in the b lattice parameter should have an impact on t' , while changes in c should affect t . In Fig. 7 (b), we see that the c/b lattice parameter ratio decreases with temperature down to $T = 150$ K, then increases with falling temperature until $T = 20$ K. This has an immediate impact on the degree of frustration t'/t as it should be approximately proportional to the c/b ratio. Indeed, comparison of Fig. 6 (b) and Fig. 7 (b) confirms this expectation and thus explains the nonmonotonous temperature dependence of the frustration.

An increasing frustration upon cooling, reaching t'/t

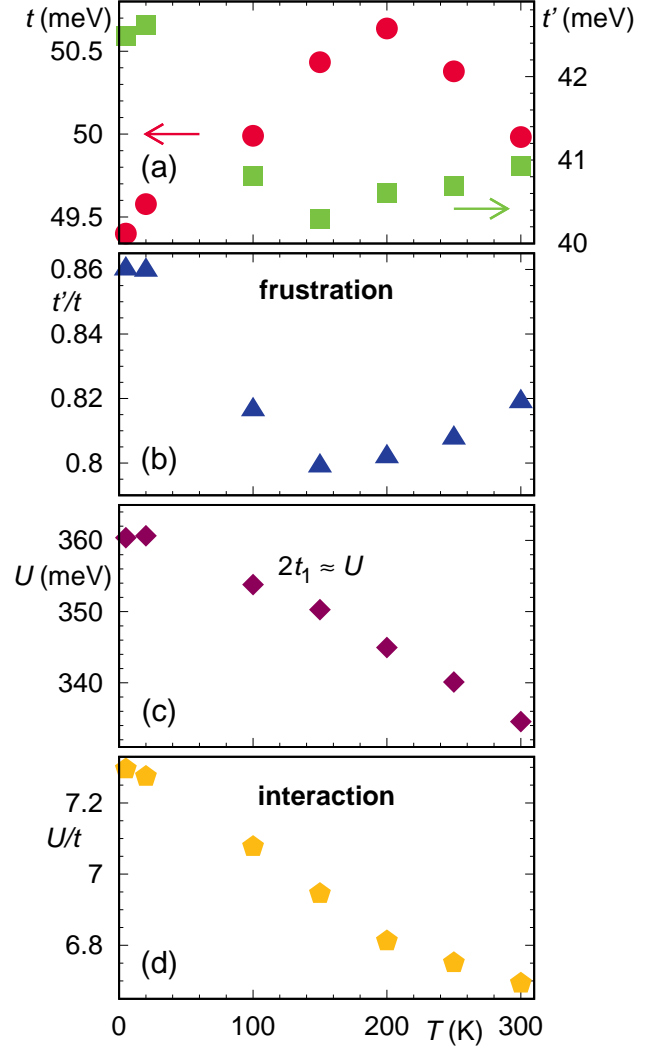


FIG. 6: Temperature dependence of Hamiltonian parameters of κ -(BEDT-TTF) $_2$ Cu $_2$ (CN) $_3$.

values at low temperatures in excess of those at room temperature, is an interesting finding which may help to better understand the intriguing low-temperature magnetic and dielectric properties of this material. Here we mention the distinct types of charge ordering, accompanied by dielectric anomalies, proposed in Ref. 17 for the present material as a result of an increasing degree of frustration.

The nonmonotonous evolution of the in-plane distortion c/b , which adopts a minimum around 150 K, might be related to the ordering of the ethylene groups, uncovered in the present study.³⁷ According to our structural analysis, the fraction of ethylene groups, ordered in the staggered conformation, gradually grows from 77% at room temperature over 86% at 250 K to 93% at 200 K. For temperatures below 150 K, the ordering is complete within the accuracy of our analysis/refinement. The large step size of 50 Kelvin employed in this study does not allow to determine the ordering temperature more pre-

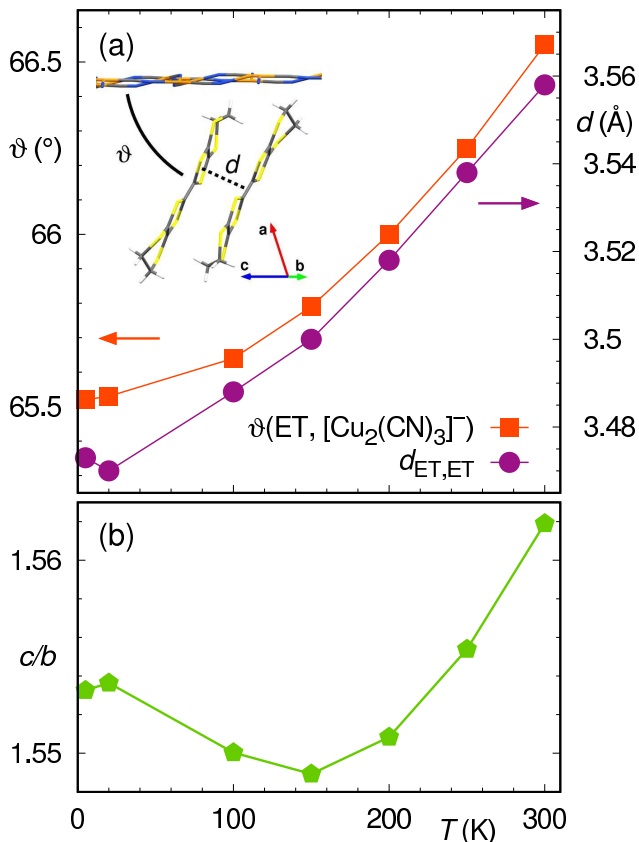


FIG. 7: (a) Orientation ϑ and dimerization d of BEDT-TTF molecules as function of temperature. The inset illustrates how ϑ and d are measured. (b) Ratio c/b between c and b lattice vectors as function of temperature.

cisely. Likewise, we cannot say whether or not the ordering occurs continuously or sets in abruptly. An argument in favour of the latter possibility might be derived from a small step-like feature revealed in thermal expansion measurements around 150 K, see the out-of plane data shown in Fig. 1 of Ref. 9. We stress, that a fully ordered staggered ethylene conformation below 150 K and a progressively disordered state above 200 K, is fully consistent with the anomalous behavior revealed by ^1H -NMR measurements between 150-200 K.⁶ The strong increase in $(T_1T)^{-1}$ above 200 K was attributed to thermally activated vibrations of ethylene groups.⁶ We suggest that the nonmonotonous temperature dependence in t'/t might

also be related to the drastic change in the thermopower around 150 K.⁷ The thermopower is related to the energy derivative of the density of states at the Fermi level.³⁸ However, for a strongly correlated system like κ -(BEDT-TTF)₂Cu₂(CN)₃, density functional theory is not sufficient for the calculation of this quantity and more elaborate many-body calculations – which are beyond the scope of the present study – have to be done.

In summary, we performed an analysis of the temperature dependence of the structural and electronic properties of κ -(BEDT-TTF)₂Cu₂(CN)₃ by considering a combination of X-ray diffraction at various temperatures and density functional calculations. Our study shows that the temperature dependence of the structural parameters has significant influence on the electronic properties and results in a nonmonotonous behavior of the degree of frustration. Of special relevance is the increase of frustration at low temperatures in comparison to the behavior at room temperature.

Acknowledgments

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Appendix A: Additional crystallographic material

Crystallographic data for the κ -(BEDT-TTF)₂Cu₂(CN)₃ structure at 5, 10, 100, 150, 200, 250 and 300 K has been deposited with the Cambridge Crystallographic Data Centre as supplementary publication nos. CCDC 850022 to 850028. Copies of the data can be obtained free of charge on application to CCDC, 12 Union Road, Cambridge CB2 1EZ, UK (fax: (44) 1223 336-0333; e-mail: data_request@ccdc.cam.ac.uk).

Fig. 8 illustrates geometry of the anion layer and the slight deviation from planarity. Fig. 9 and Table II contain lattice parameters determined at 10 K intervals between 290 K and 100 K.

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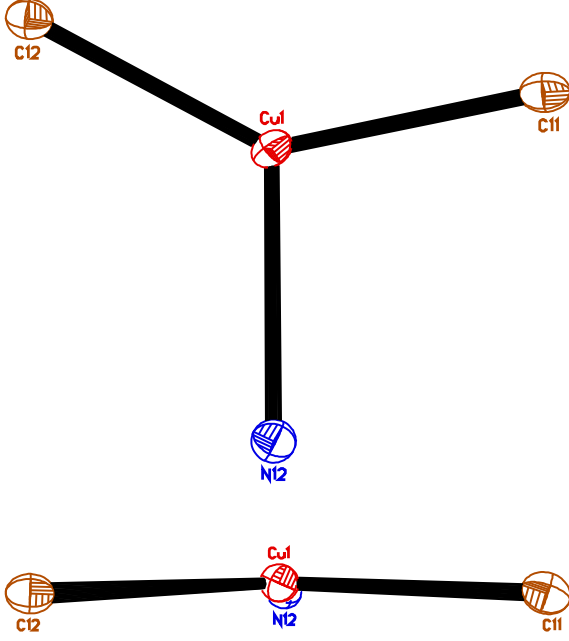


FIG. 8: The CuI coordination sphere at 5 K. The CuI is 0.072 Å out of the plane defined by N12, C11 and C12. Thermal ellipsoids are drawn at the 50% probability level.

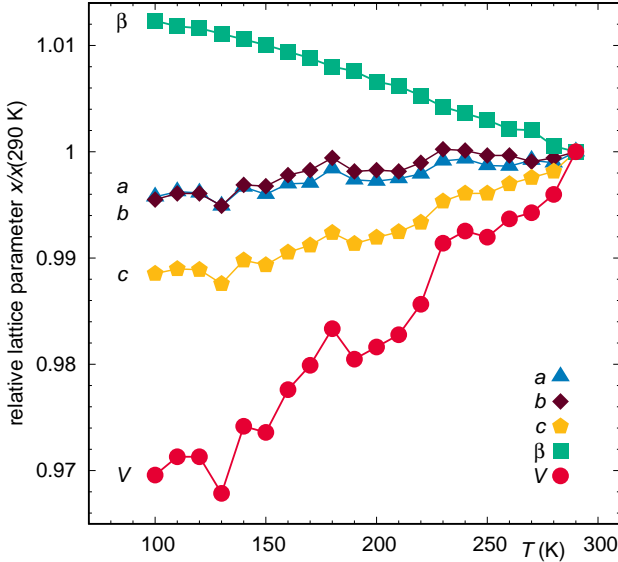


FIG. 9: Plot of the unit cell parameters as a function of temperature between 100 and 290 K.

T (K)	a (Å)	b (Å)	c (Å)	β (°)	V (Å ³)
290	16.248	8.655	13.516	113.56	1742
280	16.231	8.650	13.491	113.62	1735
270	16.235	8.647	13.483	113.79	1732
260	16.226	8.652	13.475	113.80	1731
250	16.227	8.652	13.463	113.90	1728
240	16.237	8.656	13.463	113.97	1729
230	16.234	8.657	13.453	114.04	1727
220	16.214	8.646	13.426	114.16	1717
210	16.207	8.639	13.414	114.26	1712
200	16.203	8.640	13.407	114.31	1710
190	16.205	8.639	13.399	114.42	1708
180	16.222	8.650	13.413	114.47	1713
170	16.200	8.640	13.397	114.56	1707
160	16.199	8.636	13.388	114.63	1703
150	16.183	8.627	13.372	114.70	1696
140	16.194	8.628	13.378	114.76	1697
130	16.165	8.611	13.348	114.82	1686
120	16.185	8.621	13.366	114.88	1692
110	16.187	8.621	13.367	114.90	1692
100	16.179	8.616	13.361	114.96	1689

TABLE II: Unit cell parameters as a function of temperature in the temperature range from 290 K down to 100 K.

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